

Fibrancy of Symplectic Homology in Cotangent Bundles

Thomas Kragh

April 5, 2013

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- ▶ Let N be a closed smooth manifold, then T^*N has a canonical 1-form λ defined by

$$\lambda_{(q,p)}(v) = p(\pi_*(v)),$$

where $q \in N$, $p \in T_q^*N$, $v \in T_{(q,p)}(T^*N)$ and $\pi: T^*N \rightarrow N$.

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- ▶ Ex: (DT^*N, λ) is a Liouville domain - given any Riemannian structure on N .

Exact Liouville sub-domains

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- ▶ If j is an exact Lagrangian embedding then the extension is exact.

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- ▶ Fukaya, Seidel and Smith's result was proven independently using slightly different techniques by Nadler.

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- ▶ The critical points of A_H are given precisely by the 1-periodic orbits of the Hamiltonian flow of H .
- ▶ Recall that the Hamiltonian flow is defined as the flow of X_H where X_H , solves $\omega(X_H, -) = dH$.

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- ▶ Then one defines the Floer homology $FH_*(H)$ as the Morse homology of A_H given by

$$FC_*(H) = (\mathbb{Z}[\text{critical points of } A_H], \partial),$$

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- ▶ For a a regular value we can restrict to $A_H^{-1}([a, \infty))$. We denote the resulting complex $FC_*^a(H)$.
- ▶ There is a natural quotient map

$$FH_*^a(H) \rightarrow FH_*^b(H) \quad \text{for} \quad a < b$$

Collar Neighborhood and action

- ▶ We may find a collar

$$X = (1 - \varepsilon, 1] \times \partial M \subset M$$

such that $\lambda' = u\lambda'_{|\partial M}$ on X for $u \in (1 - \varepsilon, 1]$.

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- ▶ Assume $H: X \rightarrow \mathbb{R}$ is given by $H(u, x) = h(u)$ then we get that the Hamiltonian flow preserves each leaf $\{u\} \times \partial M$, and that the action of a 1-periodic orbit is given by

$$A_H(\gamma) = \int_{\gamma} \lambda - H dt = \int_{\gamma} \lambda' - H(dt) = uh'(u) - h(u),$$

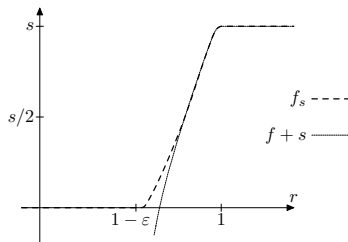
where $\{u\} \times \partial M$ is the leaf it lies in.

Symplectic Homology

- ▶ Now, fix a smooth map $f: (1 - \varepsilon, \infty) \rightarrow \mathbb{R}$ such that
 - ▶ f is concave,
 - ▶ $f(t) = 0$ for $t > 1$, and
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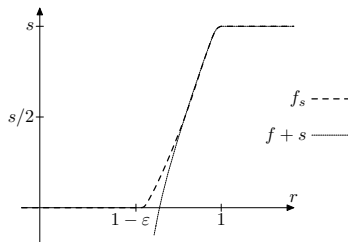
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- ▶ Now define $f_s: \mathbb{R} \rightarrow \mathbb{R}$ as

- ▶ $f_s(u) = f(u) + s$ for
 $f(u) > -s/2$.

- ▶ $f_s(u) = 0$ for $s < 1 - \varepsilon$.



- ▶ All tangents to f_s intersecting the 2. axis above 0 is tangent at points where $f_s = f + s$.

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- ▶ We now define

$$H_s(z) = \begin{cases} f_s(u) & z = (u, x) \in X \\ 0 & z \in M - X \\ s & z \in T^*N - M \end{cases}$$

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$$SH_*(M) = \operatorname{colim}_{s \rightarrow \infty} FH_*^a(H_s) \quad \text{for any } a < 0.$$

Case of $M = DT^*L$

- ▶ When $M = DT^*L$ then the flow of these Hamiltonians reproduce geodesic flows on L and Viterbo calculated that

$$SH_*(DT^*L) \cong H_*(\mathcal{L}L)$$

when L is orientable and spinable.

Fiber-wise Symplectic Homology

- ▶ We now fix $q \in N$ and look at the same action on a different space:

$$\Omega_q T^*N = \{\gamma: I \rightarrow T^*N \mid \pi(\gamma(0)) = \pi(\gamma(1)) = q\}.$$

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- ▶ It is also well known that critical points of the action defined on $\Omega_q T^*N$ are: time 1 flow curves for the Hamiltonian flow starting and ending in the fiber.
- ▶ We denote this Floer homology by $FH_*^a(H, q)$ for a regular value a .

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- ▶ Since we are dealing with paths and not loops we get an action “distortion”:

$$A_H(\gamma) = \int_{\gamma} \lambda - Hdt = \int_{\gamma} \lambda' - Hdt + f(\gamma(0)) - f(\gamma(1)),$$

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- ▶ So for a critical point γ (for A_{H_s} defined on $\Omega_q T^*N$) with critical value less than $-K$ we know precisely as before that increasing s similarly decreases the critical value, and we may thus define

$$SH_*(M, q) = \operatorname{colim}_{s \rightarrow \infty} FH_*^a(H_s, q)$$

for any $a < -K$.

Serre Type Spectral Sequence

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Theorem (K)

There is a Serre type spectral sequence strongly converging to $SH_(M)$ with second page isomorphic to $H_*(N, SH_*(M, \bullet))$.*

Sketch of Local System Lemma

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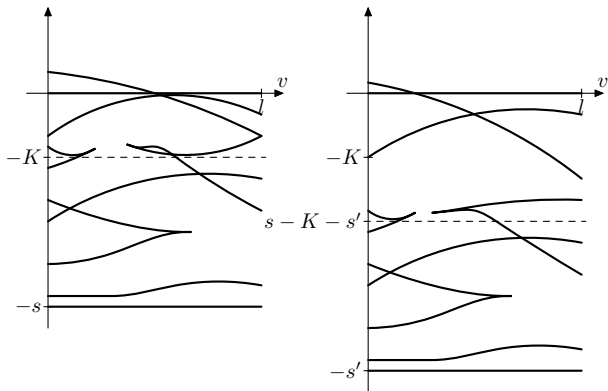
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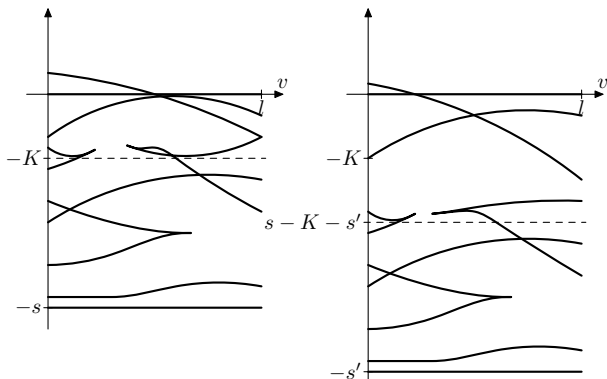
$$SH_*(M, \beta(0)) \rightarrow SH_*(M, \beta(1)).$$

- ▶ To see how, we look at the critical points of A_{H_s} restricted to $\Omega_{\beta(v)} T^*N$ for $v \in I$. This defines a so-called bifurcation diagram.

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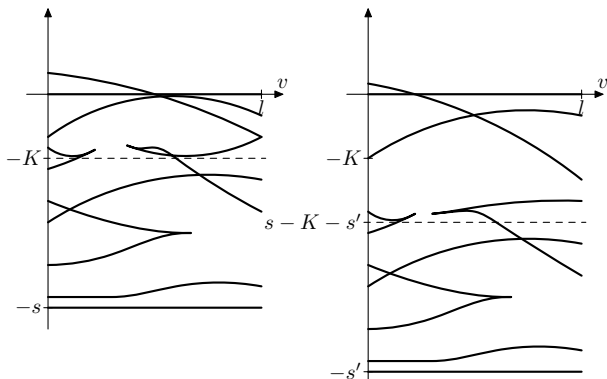


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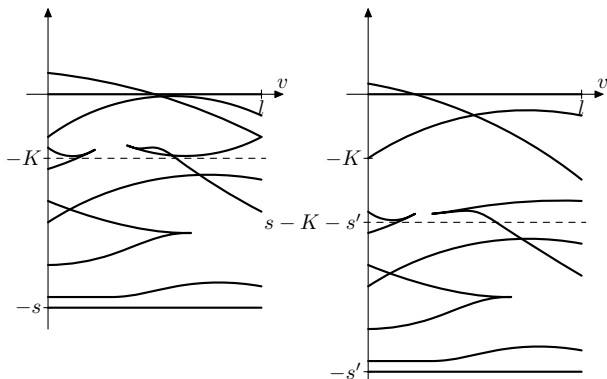
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- ▶ This implies that the similar bifurcation diagram for $A_{H_{s+Cv}}$ satisfies that all pieces below $-K$ has negative slope.

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- ▶ This means that we can (as before) produce the maps by simply collapsing generators which pass through $a < -K$. That is we produce maps

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- ▶ independent of s , and compatible with increasing s .
- ▶ Taking the colimit as $s \rightarrow \infty$ then defines the map

$$SH_*(M, \beta(0)) \rightarrow SH_*(M, \beta(1)).$$

Sketch of Local System Lemma

- ▶ This means that we can (as before) produce the maps by simply collapsing generators which pass through $a < -K$. That is we produce maps

$$FH_*^a(H_s, \beta(0)) \rightarrow FH_*^a(H_{s+C}, \beta(1))$$

- ▶ independent of s , and compatible with increasing s .
- ▶ Taking the colimit as $s \rightarrow \infty$ then defines the map

$$SH_*(M, \beta(0)) \rightarrow SH_*(M, \beta(1)).$$

- ▶ Finally one may prove that this map only depends on the homotopy type of β fixing the end points.

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- ▶ This usually constructs spaces (instead of the chain complexes), but as the Hamiltonians get more and more complicated the dimension increases, and in fact what one gets are spectra.
- ▶ So I construct a fibration of spectra, and these naturally have a serre spectral sequence.